

Statistics of limit root bundles and towards F-Theory MSSMs

Martin Bies

University of Pennsylvania

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Work with M. Cvetič, R. Donagi, M. Liu, M. Ong – 2102.10115, 2104.08297

What is this talk is about?

My goal: Answer one of the biggest open questions in theoretical physics

Find minimally supersymmetric standard model (MSSMs) solutions to string theory.

→ Study high-dimensional geometries subject to physically inspired conditions.

Approaches and (technical) challenges

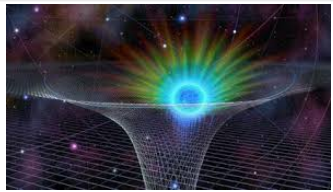
- Toric geometry: Model **many** geometries and investigate with **computers**
- Tools clustered over many different systems: Calabi-Yau tools, cohomCAlg, homalg project, Macaulay2, Sage, Singular, Topcom, Palp, Polymake, ...

⇒ My vision: OSCAR has **ALL** functionality (for F-theory studies).

Content of today's talk

- Explain state-of-the-art of my big goal.
- Explain the “**ALL**” in my vision for toric and algebraic geometry in OSCAR.

String theory = General relativity + Standard Model?



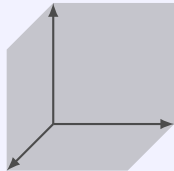
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String theory = General relativity + Standard Model?

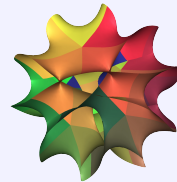


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our 4-dim. world \mathcal{W}

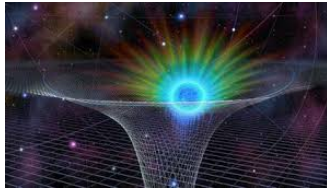
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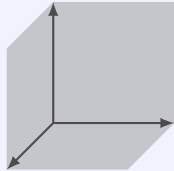
'small' 3 complex-dim. manifold \mathcal{B}_3

Challenge: Find \mathcal{B}_3 s.t. ST reproduces **Standard model**

String theory = General relativity + Standard Model?

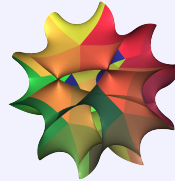


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our 4-dim. world \mathcal{W}

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'small' 3 complex-dim. manifold \mathcal{B}_3

Challenge: Find \mathcal{B}_3 s.t. ST reproduces

Particles in the standard model

Three Generations of Matter (Fermions)

	I	II	III	
mass→	2.4 MeV	1.27 GeV	171.2 GeV	0
charge→	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0
spin→	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1
name→	u up	c charm	t top	Y photon
	4.8 MeV	104 MeV	4.2 GeV	0
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	0
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1
Quarks	d down	s strange	b bottom	g gluon
	<2.2 eV	<0.17 MeV	<15.5 MeV	91.2 GeV
	0	$\frac{1}{2}$	$\frac{1}{2}$	0
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	Z weak force
	0.511 MeV	105.7 MeV	1.777 GeV	80.4 GeV
	-1	-1	-1	± 1
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1
Leptons	e electron	μ muon	τ tau	W weak force

Bosons (Forces)

- Each particle given by representation of group

$$G_{SM} = SU(3)_C \times SU(2)_L \times U(1)_Y .$$

- **Leptons** e, μ, τ : $(1, 2)_{-1/2}$
- Charge conjugate leptons e_C, μ_C, τ_C : $(\bar{1}, \bar{2})_{+1/2}$
- (e, e_C) is leptonic vector-like pair – **not** observed!
- Famous vector-like pair: Higgs field.

Necessary criterion for String theory MSSM

Number of vector-like pairs must match the experimental findings.

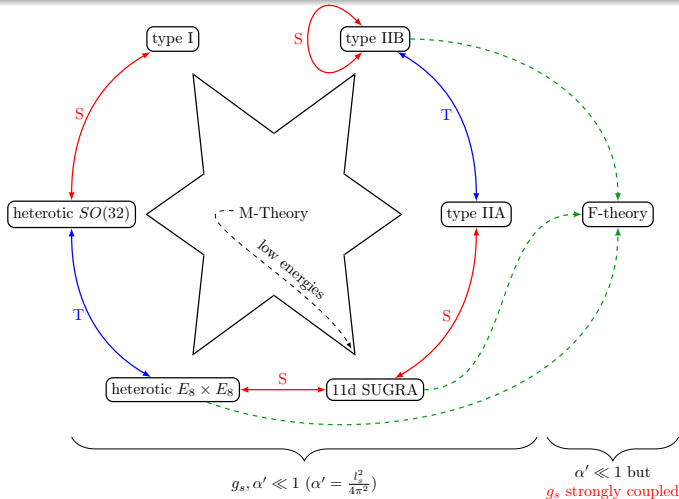
Experimentally observed vector-like spectra (LH/RH is for left/right handed)

Particle name	Rep. R	# Fields n_R	Conjugate rep. \bar{R}	# Fields $n_{\bar{R}}$
LH quarks	$(3, 2)_{1/6}$	3	$(\bar{3}, \bar{2})_{-1/6}$	0
RH up-quarks	$(3, 1)_{2/3}$	3	$(\bar{3}, \bar{1})_{-2/3}$	0
RH down-quarks	$(3, 1)_{-1/3}$	3	$(\bar{3}, \bar{1})_{1/3}$	0
LH leptons	$(1, 2)_{-1/2}$	3	$(\bar{1}, \bar{2})_{1/2}$	0
neutrinos	$(1, 1)_1$	3	$(\bar{1}, \bar{1})_{-1}$	0
Higgs	$(1, 2)_{-1/2}$	1	$(\bar{1}, \bar{2})_{1/2}$	1

Note

- Chiral index $n_R - n_{\bar{R}}$ is topological invariant – used as simple criterion.
- The “magic number” 3 follows from experiments. (→ Explain with string theory?)
- Crucial: We need **one** vector-like pair to accommodate the Higgs field.

Different formulations of string theory – the M-theory star



Overview of SM constructions in string theory

Obtain (MS)SM from String theory construction ...

- $E_8 \times E_8$ [Candelas Horowitz Strominger Witten '85], [Greene Kirklín Miron Ross '86], [Braun He Ovrut Pantev '05], [Bouchard Donagi '05], [Anderson Gray He Lukas '10], ...
- Type II [Berkooz Douglas Leigh '96], [Aldazabal Franco Ibanez Rabadan Uranga '00], [Ibanez Marchesano Rabadan '00], [Blumenhagen Kors Lust Ott '01], [Cvetič Shiu Uranga '01],
- F-theory [Krause Mayrhofer Weigand '12], [Cvetič Klevers Pena Oehlmann Reuter '15], [Lin Weigand '16], [Cvetič Lin Liu Oehlmann '18], [Cvetič Halverson Lin Liu Tian '19], ...

... including vector-like spectra

- $E_8 \times E_8$: [Bouchard Donagi '05], [Bouchard Cvetič Donagi '06], [Anderson Gray Lukas Palti '10 & '11], ...
- F-theory: [M.B. Mayrhofer Pehle Weigand '14], [M.B. Mayrhofer Weigand '17], [M.B. '18], [M.B. Cvetič Donagi Lin Liu Ruehle '20], [M.B. Cvetič Donagi Liu Ong '21], [M.B. Cvetič Liu '21]

Why F-theory

1 Conceptual reasons:

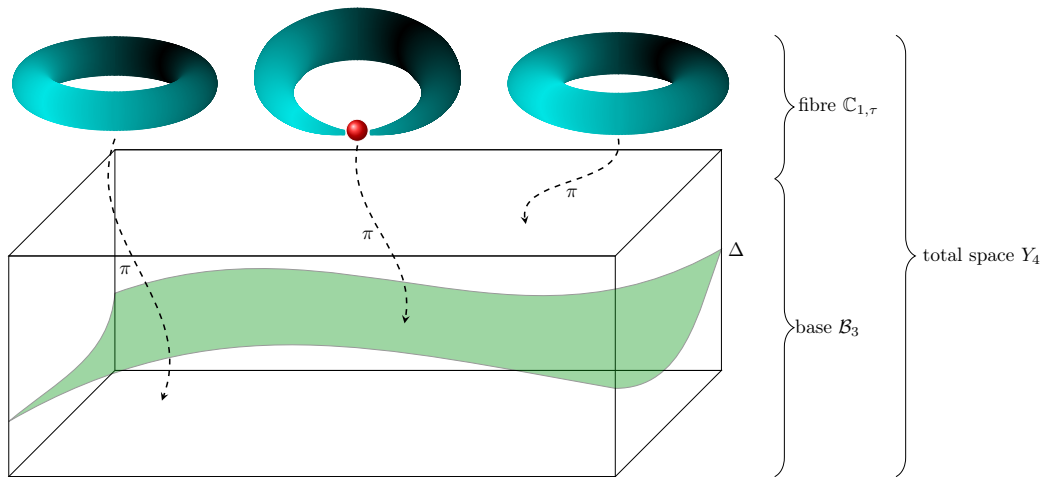
[Vafa '96], [Morrison Vafa '96], [Beasley Heckman Vafa '08], [Apruzzi Heckman Morrison Tizzano '18], ...

- More general physics setups (e.g. more general gauge groups).
- Access to strongly-coupled regime of string theory.

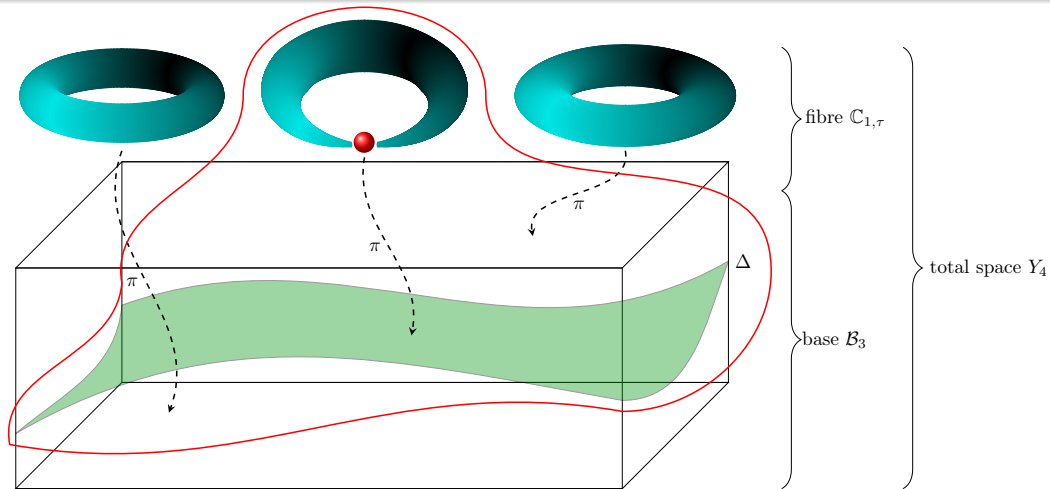
2 Model building:

- Consistency conditions: Geometry of **elliptic fibration** [Vafa '96], [Morrison Vafa '96], ...
- Yukawa couplings: **Matter curve** intersection [Cecotti Cheng Heckman Vafa '10], [Donagi, Wijnholt '12], [Cvetic Lin Liu Zhang Zoccarato '19], ...
- **Largest** (currently-known) family of string theory standard models from F-theory:
 - *A Quadrillion (10^{15}) F-theory Standard Models* [Cvetič Halverson Lin Liu Tian '19]
 - Involve toric 3-folds from triangulations of reflexive polytopes [Kreuzer Skarke '98].

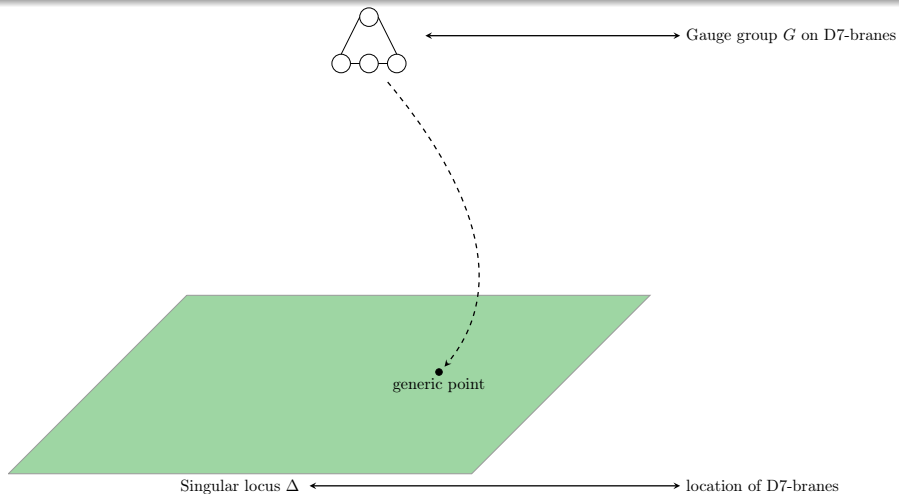
F-theory from singular fibration $Y_4 \rightarrow B_3$



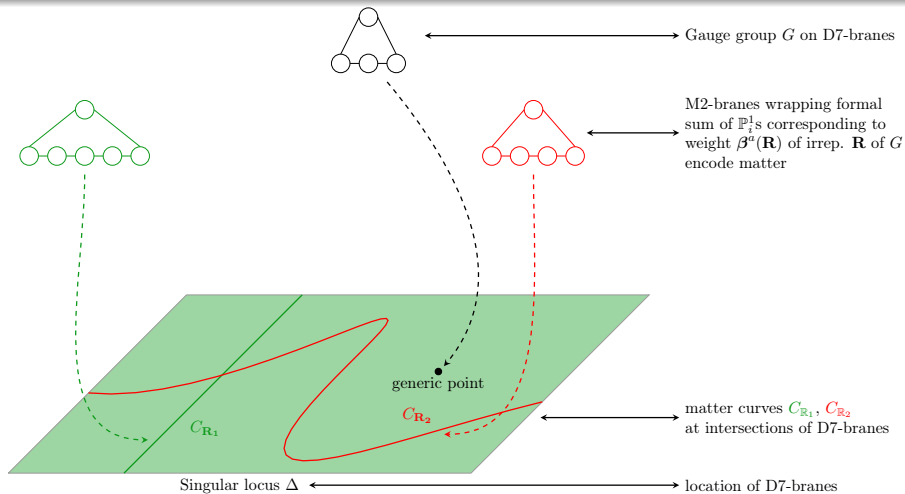
F-theory from singular fibration $Y_4 \rightarrow B_3$



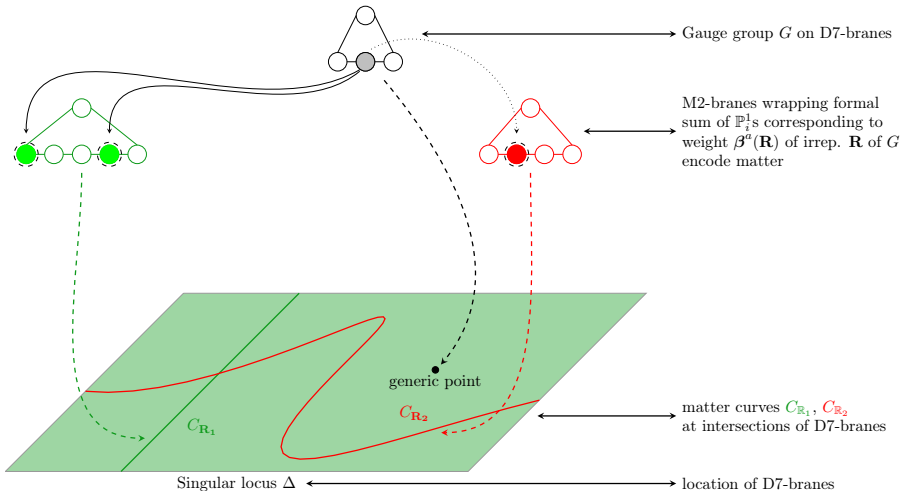
F-theory from resolved smooth fibration $\widehat{Y}_4 \rightarrow B_3$



F-theory from resolved smooth fibration $\widehat{Y}_4 \rightarrow B_3$



F-theory from resolved smooth fibration $\widehat{Y}_4 \rightarrow B_3$



Vector-like spectra in 4d $\mathcal{N} = 1$ F-theory vacua

[M.B. Mayrhofer Pehle Weigand '14], [M.B. Mayrhofer Weigand '17], [M.B. '18]

- Defining data: Elliptic 4-fold $\widehat{Y}_4 \rightarrow B_3$ and flux $G_4 \in H_{\mathbb{Z}}^{(2,2)}(\widehat{Y}_4)$.
- ⇒ Can “read-off” the physics [Weigand '17]
 - Particles localize on **matter curves** $C_R \subset B_3$.
 - Representation encoded by **matter surface** S_R (a \mathbb{P}^1 -fibration over C_R).
 - G_4 and S_R specify line bundle L_R on C_R (details on next slide):
 - massless chiral supermultiplets in rep. $R \leftrightarrow n_R = h^0(C_R, L_R)$,
 - massless chiral supermultiplets in rep. $\bar{R} \leftrightarrow n_{\bar{R}} = h^1(C_R, L_R)$.
- Challenges:
 - $n_R, n_{\bar{R}}$ **strongly** depend on the complex structure of \widehat{Y}_4 .
 - Deformation $\widehat{Y}_4 \rightarrow \widehat{Y}'_4$ can lead to **jumps** [M.B. Cvetič Donagi Lin Liu Ruehle '20]

$$h^i(C_R, L_R) = (h^0, h^1) \rightarrow h^i(C'_R, L'_R) = (h^0 + \mathbf{a}, h^1 + \mathbf{a}).$$
 - Higgs field: Need **non-generic solution** $h^i(C_R, L_R) = (1, \mathbf{1})$.

How to compute L_R from G_4, S_R ? [M.B. Mayrhofer Pehle Weigand '14], [M.B. Mayrhofer Weigand '17], [M.B. '18]

- Lift $G_4 \in H_{\mathbb{Z}}^{(2,2)}(\widehat{Y}_4)$ to a “gauge field” $A \in H_D^4(\widehat{Y}_4, \mathbb{Z}(2))$ or $\mathcal{A} \in \text{CH}^2(\widehat{Y}_4, \mathbb{Z})$:

$$\begin{array}{ccccccc}
 0 & \longrightarrow & \text{CH}_{\text{hom}}^2(\widehat{Y}_4, \mathbb{Z}) & \longrightarrow & \text{CH}^2(\widehat{Y}_4, \mathbb{Z}) & \xrightarrow{\gamma} & H_{\text{alg}}^{(2,2)}(\widehat{Y}_4) \longrightarrow 0 \\
 & & \downarrow & & \downarrow \widehat{\gamma} & & \downarrow \\
 0 & \longrightarrow & J^2(\widehat{Y}_4) & \longrightarrow & H_D^4(\widehat{Y}_4, \mathbb{Z}(2)) & \xrightarrow{\widehat{c}} & H_{\mathbb{Z}}^{(2,2)}(\widehat{Y}_4) \longrightarrow 0
 \end{array}$$

Always exists, but is in general **not unique** since $J^2(\widehat{Y}_4) \neq 0$.

- For matter surface $S_R \in \text{CH}^2(\widehat{Y}_4, \mathbb{Z})$ define $\iota_{S_R}: S_R \hookrightarrow \widehat{Y}_4$, $\pi_{S_R}: S_R \twoheadrightarrow C_R$. Then

$$L_R(\mathcal{A}) = \mathcal{O}_{C_R} [\pi_{S_R*} (\iota_{S_R}^* (\mathcal{A})) + D_{\text{spin},R}] \in \text{Pic}(C_R).$$

Summary: Computing vector-like spectra in F-theory

- 1 Defining data: Elliptic 4-fold $\widehat{Y}_4 \rightarrow B_3$ and flux $G_4 \in H_{\mathbb{Z}}^{(2,2)}(\widehat{Y}_4)$.
- 2 Read-off matter curves $C_R \subset B_3$.
- 3 Compute line bundles $L_R \in \text{Pic}(C_R)$ induced from G_4 .
- 4 Compute cohomologies to determine vector-like spectra:

Particle	R	n_R	desired $(n_R, n_{\bar{R}})$
LH quarks	$(3, 2)_{1/6}$	$h^0(C_{(3,2)_{1/6}}, L_{(3,2)_{1/6}})$	(3,0)
RH up-quarks	$(3, 1)_{2/3}$	$h^0(C_{(3,1)_{2/3}}, L_{(3,1)_{2/3}})$	(3,0)
RH down-quarks	$(3, 1)_{-1/3}$	$h^0(C_{(3,1)_{-1/3}}, L_{(3,1)_{-1/3}})$	(3,0)
LH leptons	$(1, 2)_{-1/2}$	$h^0(C_{(1,2)_{-1/2}}, L_{(1,2)_{-1/2}})$	(3,0)
neutrinos	$(1, 1)_1$	$h^0(C_{(1,1)_1}, L_{(1,1)_1})$	(3,0)
Higgs	$(1, 2)_{-1/2}$	$h^0(C_{(1,2)_{-1/2}}, L_{(1,2)_{-1/2}})$	(1, 1)

Questions?



Generalities of the QSMs

- QSMs: $\mathcal{O}(10^{15})$ elliptic 4-folds \widehat{Y}_4 with choice of G_4 [Cvetič Halverson Lin Liu Tian '19]
 - Elliptic 4-folds $\widehat{Y}_4 \twoheadrightarrow B_3$:
 - Obtained from toric geometry.
 - Constraints: no chiral exotics, massless $U(1)$ -gauge boson, cancel D_3 -tadpole.
- $\Rightarrow B_3$ from triangulations of 708 3-dim reflexive polytopes [Kreuzer Skarke '98]

$$\overline{K}_{B_3} \cdot \overline{K}_{B_3} \cdot \overline{K}_{B_3} \in \{\cancel{2}, 6, 10, 18, 30, \cancel{90}\}.$$

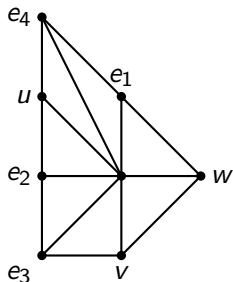
($\overline{K}_{B_3}^3 = 2, 90$ not realized by toric 3-folds.)

- Largest polytope Δ_8° : 39 lattice points and $\sim 10^{15}$ triangulations [Halverson Tian '17].
- G_4 -flux candidate: (\leftrightarrow satisfies necessary conditions to be integral)

$$G_4 = \frac{-3}{\overline{K}_{B_3}^3} \cdot (5[e_1] \wedge [e_4] + \dots) \in H_{\text{alg}}^{(2,2)}(\widehat{Y}_4).$$

Toric 5-fold ambient space for QSMs

- \widehat{Y}_4 is hypersurface in toric 5-fold space $X_5 = B_3 \times \mathbb{P}_{F_{11}}$.
- B_3 from triangulations of 708 3-dim reflexive polytopes [Kreuzer Skarke '98]
- $\mathbb{P}_{F_{11}}$ a particular toric surface [Klevers Pena Oehlmann Piragua Reuter '14]:



	u	v	w	e ₁	e ₂	e ₃	e ₄
H	1	1	1				
E ₁	-1		-1	1			
E ₂	-1	-1			1		
E ₃		-1			-1	1	
E ₄	-1			-1			1

$$I_{\text{SR}}(\mathbb{P}_{F_{11}}) = \langle e_4 w, e_4 e_2, e_4 e_3, e_4 v, e_1 u, e_1 e_2, e_1 e_3, e_1 v, wu, we_2, we_3, ve_2, uv, e_3 u \rangle.$$

\widehat{Y}_4 as hypersurface in X_5

- Consider sections $s_1, s_2, s_3, s_5, s_6, s_9 \in H^0(B_3, \overline{K}_{B_3})$.
- Define $\widehat{Y}_4 = V(p_{F_{11}}) \in X_5 = B_3 \times \mathbb{P}_{F_{11}}$ with

$$p_{F_{11}} = s_1 e_1^2 e_2^2 e_3 e_4^4 u^3 + s_2 e_1 e_2^2 e_3^2 e_4^2 u^2 v + s_3 e_2^2 e_3^3 uv^2 + s_5 e_1^2 e_2 e_4^3 u^2 w \\ + s_6 e_1 e_2 e_3 e_4 uvw + s_9 e_1 vw^2.$$

- $H^0(\mathbb{P}_{F_{11}}, \overline{K}_{\mathbb{P}_{F_{11}}})$ has basis

$$\{e_1^2 e_2^2 e_3 e_4^4 u^3, e_1 e_2^2 e_3^2 e_4^2 u^2 v, e_2^2 e_3^3 uv^2, e_1^2 e_2 e_4^3 u^2 w, e_1 e_2 e_3 e_4 uvw, e_1 vw^2\}.$$

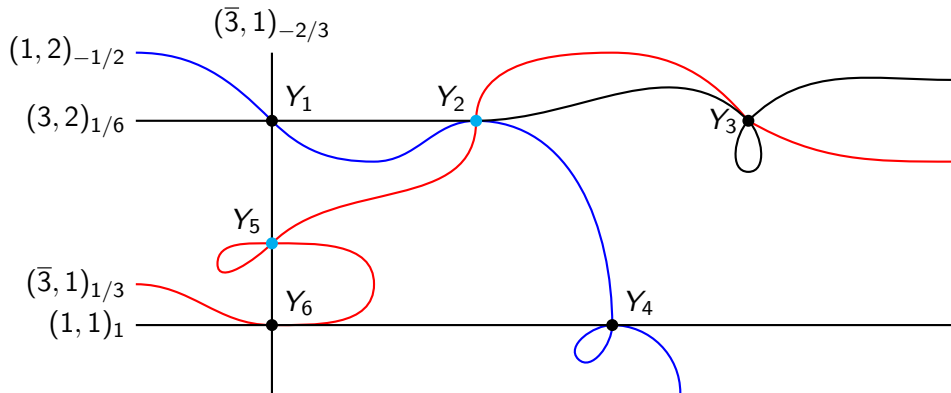
$\Rightarrow \widehat{Y}_4$ is Calabi-Yau.

Matter curves

- \widehat{Y}_4 is resolution of a singular elliptic fibration Y_4 .
- ⇒ Can read-off the physics:
 - Codimension-1 loci over which the fiber of Y_4 is singular are *gauge surfaces*.
 - Codimension-2 loci over which the fiber of Y_4 is more singular than over the gauge surfaces are *matter curves*.
- ↔ depends on choice of s_i .
- For generic s_i have five matter curves: [Klevers Pena Oehlmann Piragua Reuter '14]

$$\begin{aligned} C_{(3,2)_{1/6}} &= V(s_3, s_9), & C_{(1,2)_{-1/2}} &= V(s_3, s_2 s_5^2 + s_1(s_1 s_9 - s_5 s_6)), \\ C_{(\bar{3},1)_{-2/3}} &= V(s_5, s_9), & C_{(\bar{3},1)_{1/3}} &= V(s_9, s_3 s_5^2 + s_6(s_1 s_6 - s_2 s_5)), \\ C_{(1,1)_1} &= V(s_1, s_5). \end{aligned}$$

A cartoon of the matter curves in the QSMs



The topological intersection number is $\overline{K}_{B_3}^3$ at Y_1, Y_3, Y_4, Y_6 and $2 \cdot \overline{K}_{B_3}^3$ at Y_2, Y_5 .

Can OSCAR help?

My hope/vision: **Automatically** perform fully fledged F-theory construction

- OSCAR has **ALL** ingredients to perform a fully-fledged F-theory compactification.
- As a starting point, focus on QSMs [Klevers Pena Oehlmann Piragua Reuter '14].

Basic ingredients

- Toric varieties from triangulation of polytopes (cf. PR-848):
 - Calabi-Yau tools, Polymake [Jordan, Joswig, Kastner '18] and Topcom
 - Kreuzer-Skarke lists [Kreuzer Skarke '98], [Kreuzer Skarke '00], [Altman Gray He Jejjala Nelson '14], ...
 - QSM list (currently in GAP-4-package QSMEexplorer)
- Complete intersection subvarieties:
 - Cohomologies: cohomCalg, Calabi-Yau tools, ToricVarieties_project, ...
 - Smoothness.

Can OSCAR help? II

Chiral and vector-like spectra

- **Topological intersection** numbers of algebraic cycles.
- ⇒ Construct (vertical) G_4 -fluxes and compute chiral spectra.
- For $\mathcal{A}, S_R \in \text{CH}^2(\widehat{Y}_4, \mathbb{Z})$ compute intersection in **Chow ring**.
- ⇒ Automatically compute line bundle L_R .

Example for intersection computation in Chow ring (cf. PR-520)

- In $\mathbb{P}_{x,y,z}^2$ compute (a) self-intersection locus of algebraic cycles $C_1 = V(x)$, $C_2 = V(x)$:
- 1 Move in general position: $C_2 = V(x) \sim C'_2 = V(y)$ (since $x - y \in I_{\text{Linear Relations}}$).
 - 2 Hence $C_1 \cdot C_2 \sim V(x) \cdot V(y) = V(x, y)$.

Questions?



Vector-like spectra in the QSMs: Lifting G_4 by easier multiple

- For the QSMs, the G_4 -flux is

$$G_4 = \frac{-3}{\overline{K}_{B_3}^3} \cdot (5[e_1] \wedge [e_4] + \dots) .$$

- Naive lift $\mathcal{A} = \frac{-3}{\overline{K}_{B_3}^3} \cdot (5V(e_1, e_4) + \dots) \notin \text{CH}^2(\widehat{Y}_4, \mathbb{Z})$ since $\frac{-3 \cdot 5}{\overline{K}_{B_3}^3} \notin \mathbb{Z}$.

↔ Lack of computational control over $J^2(\widehat{Y}_4)$.

→ Can (currently) not directly write-down lift of G_4 to $\mathcal{A} \in \text{CH}^2(\widehat{Y}_4, \mathbb{Z})$.

- Circumvent ignorance:

① Consider $G'_4 = \overline{K}_{B_3}^3 \cdot G_4$ instead.

② Lift G'_4 to $\mathcal{A}' = -3 \cdot (5V(e_1, e_4) + \dots) \in \text{CH}^2(\widehat{Y}_4, \mathbb{Z})$ and find

$$D_R(\mathcal{A}') = \pi_{S_R*} (\iota_{S_R}^*(\mathcal{A}')) \in \text{Pic}(C_R) .$$

⇒ Root bundle constraint in $\text{Pic}(C_R)$: $\overline{K}_{B_3}^3 \cdot D_R(\mathcal{A}) \sim D_R(\mathcal{A}')$.

Summary of **necessary** root bundle constraints in the QSMs

curve	constraint
$C_{(3,2)_{1/6}} = V(s_3, s_9)$	$P_{(3,2)_{1/6}}^{\otimes 2\bar{K}_{B_3}^3} = K_{(3,2)_{1/6}}^{\otimes (6+\bar{K}_{B_3}^3)}$
$C_{(1,2)_{-1/2}} = V(s_3, P_H)$	$P_{(1,2)_{-1/2}}^{\otimes 2\bar{K}_{B_3}^3} = K_{(1,2)_{-1/2}}^{\otimes (4+\bar{K}_{B_3}^3)} \otimes \mathcal{O}_{C_{(1,2)_{-1/2}}}(-30 \cdot Y_1)$
$C_{(\bar{3},1)_{-2/3}} = V(s_5, s_9)$	$P_{(\bar{3},1)_{-2/3}}^{\otimes 2\bar{K}_{B_3}^3} = K_{(\bar{3},1)_{-2/3}}^{\otimes (6+\bar{K}_{B_3}^3)}$
$C_{(\bar{3},1)_{1/3}} = V(s_9, P_R)$	$P_{(\bar{3},1)_{1/3}}^{\otimes 2\bar{K}_{B_3}^3} = K_{(\bar{3},1)_{1/3}}^{\otimes (4+\bar{K}_{B_3}^3)} \otimes \mathcal{O}_{C_{(\bar{3},1)_{1/3}}}(-30 \cdot Y_3)$
$C_{(1,1)_1} = V(s_1, s_5)$	$P_{(1,1)_1}^{\otimes 2\bar{K}_{B_3}^3} = K_{(1,1)_1}^{\otimes (6+\bar{K}_{B_3}^3)}$

(P_H, P_R are complicated polynomials, Y_1, Y_3 are Yukawa points.)

Example: Absence of exotic vector-like quark-doublets

- Study QSM geometry defined by 3-fold B_3 with $\overline{K}_{B_3}^3 \in \{6, 10, 18, 30\}$.
- Pick (generic) sections $s_3, s_9 \in H^0(B_3, \overline{K}_{B_3})$.
- $C_{(3,2)_{1/6}} = V(s_3, s_9)$ and $g = \frac{\overline{K}_{B_3}^3 + 2}{2}$ ((4, 6, 10, 16) for $\overline{K}_{B_3}^3 \in \{6, 10, 18, 30\}$)
- Number of fields: $h^0(C_{(3,2)_{1/6}}, P_{(3,2)_{1/6}})$, $h^1(C_{(3,2)_{1/6}}, P_{(3,2)_{1/6}})$ where

$$P_{(3,2)_{1/6}}^{\otimes 2\overline{K}_{B_3}^3} = K_{(3,2)_{1/6}}^{\otimes (6 + \overline{K}_{B_3}^3)}.$$

\Rightarrow Necessary condition: Prove existence of $P_{(3,2)_{1/6}}$ with

$$h^0(C_{(3,2)_{1/6}}, P_{(3,2)_{1/6}}) = 3, \quad h^1(C_{(3,2)_{1/6}}, P_{(3,2)_{1/6}}) = 0. \quad (1)$$

Price of ignorance

- Which root bundles are physical, i.e. induced from $A \in H_D^4(\widehat{Y}_4, \mathbb{Z}(2))$?
(If $g > h^{21}(\widehat{Y}_4)$, then not all are physical.)
- Interesting, but also very challenging question for future work.
- For the time being: Ignore this issue.

Local bottom-up analysis

- ① For B_3 from triangulation of Δ_{40}° , I will prove existence of solution to

$$P_{(3,2)_{1/6}}^{\otimes 2\bar{K}_{B_3}^3} = K_{(3,2)_{1/6}}^{\otimes (6+\bar{K}_{B_3}^3)} \quad \text{and} \quad h^0(C_R, P_R) = 3.$$

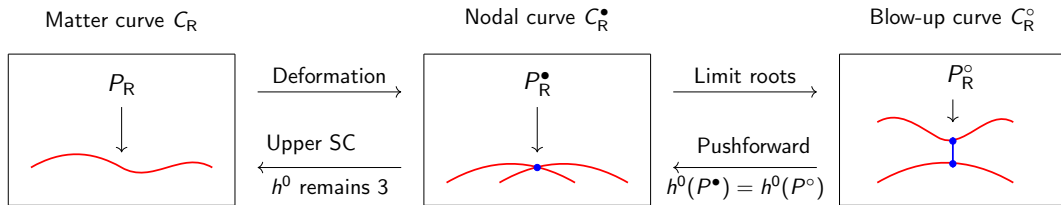
Since $\bar{K}_{B_3}^3 = 18$, it is sufficient to solve

$$P_{(3,2)_{1/6}}^{\otimes 3} = K_{(3,2)_{1/6}}^{\otimes 2} \quad \text{and} \quad h^0(C_R, P_R) = 3.$$

- ② Extend to statistical study across $\sim 10^{15}$ QSM bases.

Existence of roots from deformation theory and limit root bundles

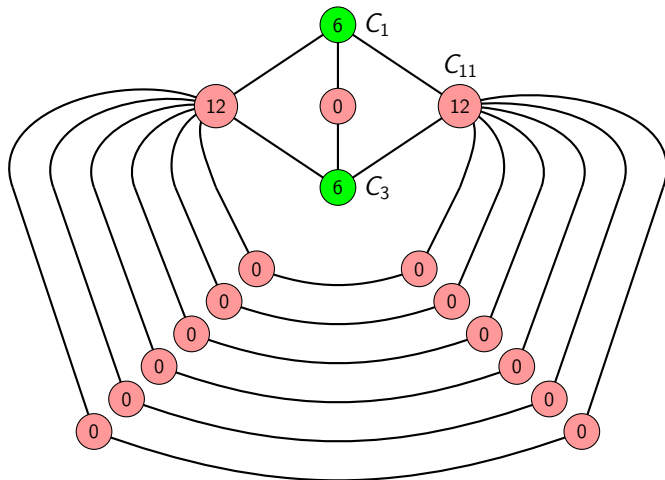
- Smooth, irreducible C_R with $g > 1$: Very hard to explicitly construct root bundles.
 - Nodal curves C_R^\bullet : well understood. [Jarvis '98], [Jarvis '99], [Caporaso Casagrande Cornalba '07]
- ⇒ Our approach is summarized as follows:



(To model *all* roots, must also consider partial blow-ups. This makes the section counting hard. Hence, we currently ignore this.)

- Deformation: $C_{(3,2)_{1/6}} = V(s_3, s_9) \rightarrow C_{(3,2)_{1/6}}^\bullet = V(s_3, \prod_i X_i)$ (\leftrightarrow Picard lattice of K3-surface)

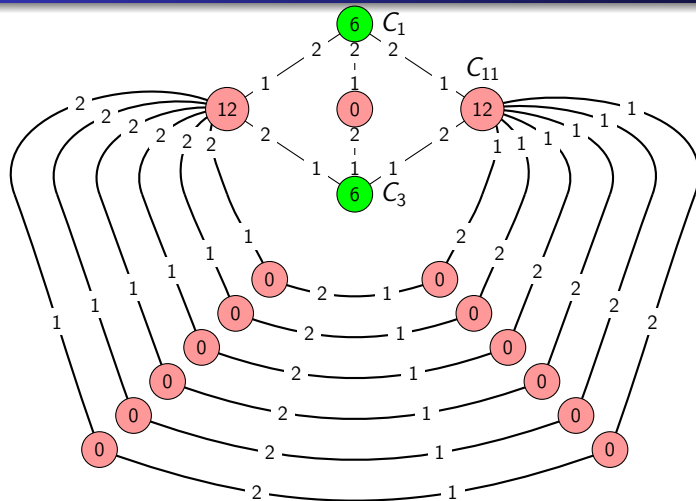
Limit root bundle construction: Step 1 – dual graph of $C_{(3,2)_{1/6}}^\bullet$



- Red bullet: $g = 0$ cpnt.
- Green bullet: $g = 1$ cpnt.
- Line: node
- Numbers: $2 \cdot \deg(K_{C_{(3,2)_{1/6}}^\bullet})$

⇒ Find 3rd roots with $h^0 = 3!$
 (Here fortunate case, as we can divide the local degrees by 3. This is not always true for QSM setups.)

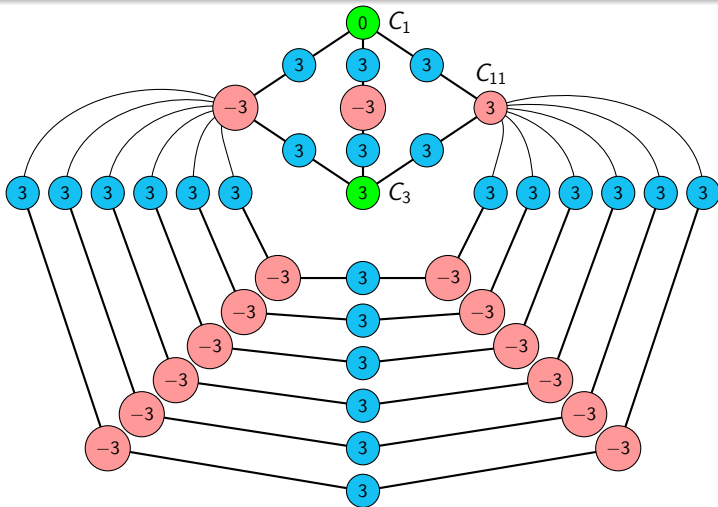
Limit root bundle construction: Step 2 – shift degrees to blow-ups $E_j \cong \mathbb{P}^1$.



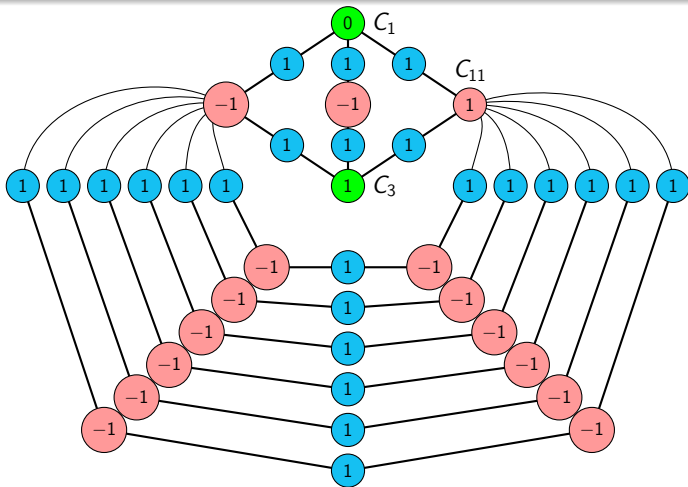
Rules for k -th roots:
 (here $k = 3$):

- $w_i \in \{1, \dots, k - 1\}$,
 - $w_1 + w_2 = k$,
 - On each component, the resulting degree is divisible by k .
- ⇒ Many possibilities!

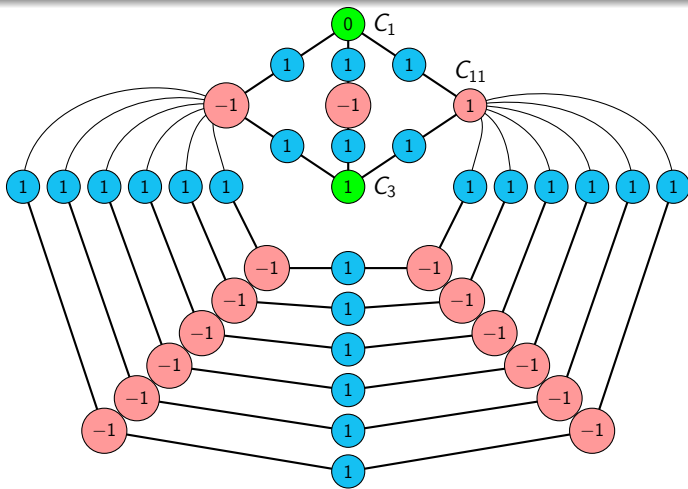
Limit root bundle construction: Step 3 – divide by $k = 3$.



Limit root bundle construction: $P_{(3,2)_{1/6}}^\bullet$ with $\left(P_{(3,2)_{1/6}}^\bullet\right)^{\otimes 3} = \left(K_{(3,2)_{1/6}}^\bullet\right)^{\otimes 2}$.



Counting h^0 of limit root bundle.



Observation:

- $h^0(E_j \cong \mathbb{P}^1, \mathcal{O}_{E_j}) = 2$
- \Rightarrow Uniquely fixed by boundary conditions.
- $\Rightarrow h^0(P_R^\bullet) = \sum_{C_i \neq E_j} h^0(C_i, P_R^\bullet|_{C_i})$
- $h^0(C_3) = 1, h^0(C_{11}) = 2$
- $h^0(C_1) = 0$ for at least 8 of 9 local roots.
- $\Rightarrow \exists P_R^\bullet$ s.t. $h^0(C_R^\bullet, P_R^\bullet) = 3.$

Towards promising F-theory base spaces

- Assume that C_R^\bullet consists only of $g = 0$ and $g = 1$ components. Then automate:
 - Given the dual graph of C_R^\bullet , find all allowed weight assignments.
 - For the encoded limit root line bundles compute $h^0(C_R^\bullet, P_R^\bullet)$.

→ Implemented in GAP-4 package *QSMExplorer*

https://github.com/homalg-project/ToricVarieties_project/tree/master/QSMExplorer

- Scan over selected QSM geometries:

① $h^{(2,1)}(\widehat{Y}_4) \geq g(C_{(3,2)_{1/6}})$.

Necessary condition for many roots on $C_{(3,2)_{1/6}}$ to be physical.

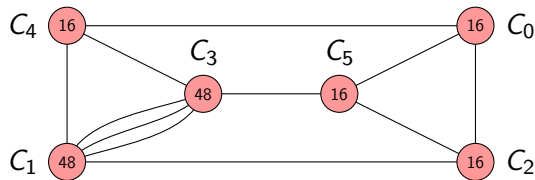
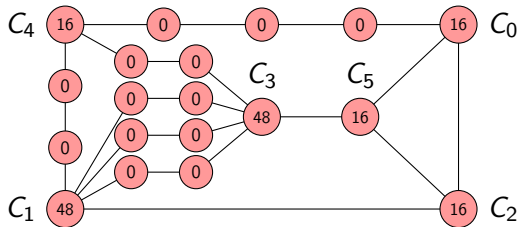
② Components of $C_{(3,2)_{1/6}}$ must have at most $g = 1$.

⇒ Base 3-folds B_3 obtained from triangulations of 33 3-dim. reflexive polytopes:

- All $\overline{K}_{B_3}^3 = 6$ bases (7 polytopes),
- Some $\overline{K}_{B_3}^3 = 10$ bases (26 polytopes).

Crucial observations which facilitate scan over $\sim 10^{15}$ different 3-folds B_3

- 1 The number of limit root bundles on C_R° with $h^0 = 3$ is independent of the triangulation! [Batyrev '93] [Cox Katz '99] [Kreuzer '06]
 - 2 Remove all tree-like subgraphs
- ⇒ Dual graph simplifies, number of limit root bundles and cohomologies unchanged!
- Example: Δ_{52}° ($\bar{K}_{B_3}^3 = 10$):



Results for bases with $\overline{K}_{B_3}^3 = 6$

- $N_P = 12^8 = (2\overline{K}_{B_3})^{2g}$: total number of root bundles on $C_{(3,2)_{1/6}}$
- $\check{N}_P^{(3)}$: number of limit roots on $C_{(3,2)_{1/6}}^\circ$ with $h^0 = 3$
- Computer scan finds:

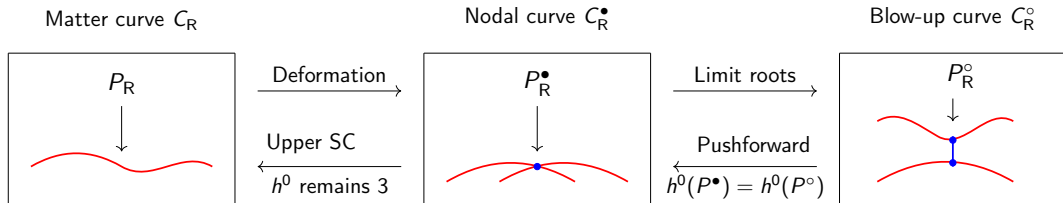
	$\check{N}_P^{(3)}$	$N_P/\check{N}_P^{(3)}$		$\check{N}_P^{(3)}$	$N_P/\check{N}_P^{(3)}$
Δ_8°	142560	$3.0 \cdot 10^3$		Δ_{130}°	$4.8 \cdot 10^4$
Δ_4°	11110	$3.8 \cdot 10^4$		Δ_{136}°	$4.8 \cdot 10^4$
Δ_{134}°	10100	$4.3 \cdot 10^4$		Δ_{236}°	$4.8 \cdot 10^4$
Δ_{128}°	8910	$4.8 \cdot 10^4$			

- Our scan is limited to a **subset** of all roots on $C_{(3,2)_{1/6}}^\bullet$.
 - $\check{N}_P^{(3)}(C_{(3,2)_{1/6}}) \geq \check{N}_P^{(3)}(C_{(3,2)_{1/6}}^\bullet)$ due to jumps along $C_{(3,2)_{1/6}}^\bullet \rightarrow C_{(3,2)_{1/6}}$.
- ⇒ Current techniques: $B_3(\Delta_8^\circ)$ most promising for F-theory MSSM (with $\overline{K}_{B_3}^3 = 6$).

Summary

- Root bundles arise naturally in the QSMs. [Cvetič Halverson Lin Liu Tian '19]
- On smooth, irreducible C_R hard, but easy for nodal C_R^\bullet :

[Jarvis '98], [Jarvis '99], [Caporaso Casagrande Cornalba '07]



- Find roots on $C_{(3,2)_{1/6}}$, $C_{(\bar{3},1)_{-2/3}}$, $C_{(\bar{3},1)_{1/3}}$, $C_{(1,1)_1}$ without vector-like exotics.
- Extend systematically to all $\mathcal{O}(10^{15})$ QSM spaces.
- With current techniques: Absence of vector-like exotics on $C_{(3,2)_{1/6}}$ most likely for base 3-folds from triangulations of Δ_8° .

Theory outlook

- Goal: F-theory MSSM construction
- ⇒ Identify roots on Higgs curve with $h^i = (4, 1)$, i.e. exactly one Higgs pair.
- Technical extensions:
 - Perform limit root counting on Higgs curve.
 - Extend limit root counting beyond limit roots on full-blowup curve C_R° .
- Conceptual questions/obstructions:
 - Does the topology of the dual graph encode the root bundle distribution?
 - ⇒ Possibly a machine learning/data science project for the summer.
 - What conditions prevent/detect jumps in vector-like spectrum along $C_R^\bullet \rightarrow C_R$?
 - More ambitious: What is the defect, i.e. by how much does h^0 jump?
 - Which root bundles are realized top-down, i.e. from an F-theory gauge potential?

Outlook for toric (and algebraic) geometry in OSCAR

- My hope/vision: **Automatically** perform fully fledged F-theory construction.
- As a starting point, focus on QSMs [Klevers Pena Oehlmann Piragua Reuter '14].
- Toric varieties from triangulation of polytopes (cf. PR-848):
 - Calabi-Yau tools, Polymake [Jordan, Joswig, Kastner '18] and Topcom
 - Kreuzer-Skarke lists [Kreuzer Skarke '98], [Kreuzer Skarke '00], [Altman Gray He Jejjala Nelson '14], ...
 - QSM list (currently in GAP-4-package QSMExplorer)
- Complete intersection subvarieties:
 - Cohomologies: cohomCAlg, Calabi-Yau tools, ToricVarieties_project, ...
 - Smoothness.
- Chiral and vector-like spectra:
 - **Topological intersection** numbers of algebraic cycles.
 - **Chow ring and intersection theory.**
- Root bundle counter (currently in GAP-4-package QSMExplorer)

Thank you for your attention!

